

# CHARGE RECOMBINATION KINETICS AS A PROBE OF PROTONATION OF THE PRIMARY ACCEPTOR IN PHOTOSYNTHETIC REACTION CENTERS

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**ABSTRACT** The kinetics of the charge recombination  $D^+Q_A^- \rightarrow DQ_A$  was used to probe the protonation of the primary acceptor in reaction centers from *Rhodospseudomonas sphaeroides*, in which the native ubiquinone was replaced by anthraquinone. We found that  $Q_A^-$  is stabilized by the rapid ( $t < 10^{-2}$  s) binding of a proton, with a pK of 9.8. The distance between  $Q_A^-$  and the proton binding site was estimated to be larger than  $\sim 5 \text{ \AA}$ .

The absorption of light by photosynthetic reaction centers (RCs) leads to a separation of charge between the electron donor, D(a bacteriochlorophyll dimer), and the primary (quinone) acceptor,  $Q_A$ , forming  $D^+Q_A^-$  (for review, see references 1 and 2). Here we examine the protonation events accompanying the charge separation. Redox titrations of the  $Q_A/Q_A^-$  couple (3–5; for review, see reference 6) indicate that  $Q_A^-$  associates with a proton ( $pK_A = 9.8$ ). The same conclusion was reached from studies on the kinetics of electron transfer between  $Q_A^-$  and the secondary acceptor,  $Q_B$  (7). Spectroscopic measurements, however, indicate that the proton is not directly bound to  $Q_A^-$ . The optical absorption spectrum of  $Q_A^-$  at neutral pH (8–10) resembles that of the unprotonated semiquinone anion (11, 12). Similarly, the EPR spectrum of  $Q_A^-$  (13–16) is characteristic of the unprotonated anion (16).

To probe for a protonation site in the vicinity of  $Q_A$ , we studied the charge recombination kinetics  $D^+Q_A^- \rightarrow DQ_A$ , with anthraquinone (AQ) replacing the native ubiquinone (UQ). With AQ, as opposed to UQ, the charge recombination rate,  $k_{\text{obs}}$ , is sensitive to small perturbations in the free energy of  $Q_A^-$  (17, 18). Thus  $k_{\text{obs}}$  should be sensitive to a shift in the energy of  $Q_A^-$  caused by the electrostatic interaction with a nearby proton. A preliminary account of this work has been presented (19).

The charge recombination process of  $D^+Q_A^-$  is schematically illustrated in Fig. 1;  $k_{\text{H}^+}^{\text{on}}$  and  $k_{\text{H}^+}^{\text{off}}$  are the rate for proton binding and release, and  $k_{\text{obs}}^{\text{H}^+}$  and  $k_{\text{obs}}^0$  are the recombination rates in the presence and absence of a proton, respectively. The electron on  $Q_A^-$  recombines with  $D^+$  via thermal repopulation of the state  $D^+I^-Q_A$ , where I

is the intermediate acceptor (2). The recombination rate depends on the free energy difference,  $\Delta G^0$ , between  $D^+I^-Q_A$  and  $D^+IQ_A^-$  (18). The shift in this energy difference,  $\delta G^0$ , caused by the binding of a proton, results in a change in the rate given by (18):

$$k_{\text{obs}}^{\text{H}^+} = k_{\text{obs}}^0 e^{-\delta G^0/k_b T}, \quad (1)$$

where  $k_b$  is Boltzmann's constant and  $T$  the absolute temperature.

The time dependence of the recombination kinetics,  $k_{\text{obs}}$ , will depend on whether the proton equilibration rate,  $k_{\text{on}}^{\text{H}^+} + k_{\text{off}}^{\text{H}^+}$  (Fig. 1), is fast or slow with respect to the charge recombination rates  $k_{\text{obs}}^{\text{H}^+}$  and  $k_{\text{obs}}^0$ . If the proton equilibration rates is fast, the states  $(D^+Q_A^-)H^+$  and  $D^+Q_A^-$  are in equilibrium on the time scale of the charge recombination; the observed kinetics will follow a single exponential decay, with  $k_{\text{obs}}$  given by the sum of  $k_{\text{obs}}^{\text{H}^+}$  and  $k_{\text{obs}}^0$  weighted by the fraction of unprotonated and protonated RCs, respectively, i.e.,

$$k_{\text{obs}} = \frac{k_{\text{obs}}^0 + 10^{(pK_A - pH)} k_{\text{obs}}^{\text{H}^+}}{1 + 10^{(pK_A - pH)}}, \quad (2)$$

where

$$(pK_A - pH) = \log \frac{[(D^+Q_A^-)H^+]}{[D^+Q_A^-]} = \log \frac{k_{\text{on}}^{\text{H}^+}}{k_{\text{off}}^{\text{H}^+}}. \quad (3)$$

When the proton equilibration rate is slow compared with  $k_{\text{obs}}^{\text{H}^+}$  and  $k_{\text{obs}}^0$ , the states  $(D^+Q_A^-)H^+$  and  $D^+Q_A^-$  do not interconvert on the time scale of the charge recombina-

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<sup>1</sup>These states contain the intermediate state, I, which we have omitted for the sake of simplicity.

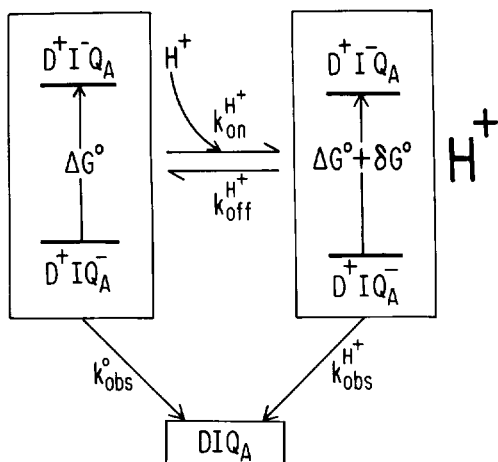


FIGURE 1 Schematic representation of the change in energy levels of  $D^+I^-Q$  and  $D^+IQ^-$  upon protonation. When the primary quinone is anthraquinone, the charge recombination  $D^+IQ^- \xrightarrow{k_{\text{obs}}} DIQ_A$  proceeds via the thermally excited state  $D^+I^-Q_A$  and depends, therefore, exponentially on the energy difference  $\Delta G^0 + \delta G^0$ . (18)

tion, and the observed kinetics will follow two exponential processes; RCs in the state  $D^+Q_A^-$  recombine with rate  $k_{\text{obs}}^0$ , while those in the state  $(D^+Q_A^-)H^+$  recombine with rate  $k_{\text{obs}}^{H^+}$ .

The recombination kinetics were measured by monitoring the optical absorption change  $\Delta A^{865}$ ; this change corresponds to the formation and subsequent decay of  $D^+$  (1). Experiments were performed using RCs isolated from *R. sphaeroides* R-26 (1), depleted of the native UQ (20), and reconstituted with either AQ or UQ (18, 20). Flash-

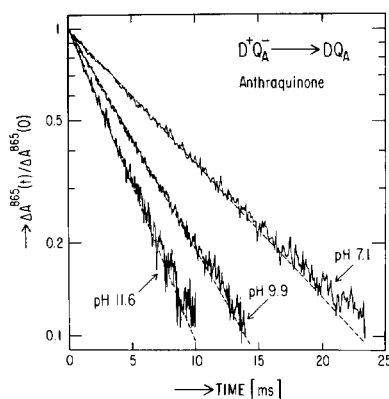


FIGURE 2 Semilogarithmic plot of the change in optical absorption at 865 nm after a single, saturating flash of light (0.4  $\mu\text{s}$  pulse width, 0.2 J/pulse) for different values of pH. The contribution to the kinetics from a residual ( $\sim 10\%$ ) population of RCs containing UQ was measured before the addition of AQ and was subtracted in the data shown. The solid lines represent the best fit of the function  $\Delta A^{865}(t)/\Delta A^{865}(0) = e^{-k_{\text{obs}}t}$  to the data. The observed kinetics were independent of the light intensity for  $10^5 \text{ W/cm}^2 < I < 10^6 \text{ W/cm}^2$ . Conditions: To 1.8  $\mu\text{M}$  RCs depleted of UQ were added, 20  $\mu\text{M}$  AQ in either 10 mM PIPES (piperazine-*N*, *N'*-bis-[2-ethanesulfonic acid]) (pH 7.1), 10 mM CHES (cyclohexylamino-ethanesulfonic acid) (pH 9.9) or 10 mM CAPS [3-(cyclohexylamino)-propanesulfonic acid] (pH 11.6), and 0.025% (wt/vol) LDAO (lauryl dimethylamine-*N*-oxide) at  $T = 21^\circ\text{C}$ .

induced charge separation was accomplished with a pulsed dye laser ( $\lambda_0 = 584 \text{ nm}$ , 0.4  $\mu\text{s}$  pulse width, 0.2 J/pulse). Changes in optical absorption were recorded with a spectrophotometer of local design (7). Experimental conditions were as described (7).

The recovery kinetics  $D^+ \rightarrow D$  are shown in Fig. 2. The decrease in the recovery rate with decreasing pH implies that  $D^+Q_A^-$  is stabilized by the presence of a proton. The kinetics followed a single-exponential decay at all pH values (Fig. 2), implying that the proton binds rapidly compared with  $k_{\text{obs}}$ , i.e.,

$$k_{\text{on}}^{H^+} + k_{\text{off}}^{H^+} \gg k_{\text{obs}} \sim 10^2 \text{ s}^{-1}. \quad (4)$$

This limit is consistent with the results of proton uptake measurements (21–23); the stoichiometry of the measured uptake is, however, controversial (21–24).

The pH dependence of  $k_{\text{obs}}$  is shown in Fig. 3. The value of  $k_{\text{obs}}$  at low pH is in agreement with previous findings (17, 18, 25). To verify that the increase in  $k_{\text{obs}}$  at high pH was not an artifact of the quinone removal and reconstitution process, measurements were performed with RCs reconstituted with UQ. The rate  $k_{\text{obs}}$  changed by  $< 25\%$  over the range  $6 < \text{pH} < 11$ , in agreement with results for native RCs (7). The solid line in Fig. 3 represents the best fit of the kinetics model (Eq. 2) to the data with  $k_{\text{obs}}^{H^+} = 97 \text{ s}^{-1}$ ,  $k_{\text{obs}}^0 = 230 \text{ s}^{-1}$  and  $\text{pK}_A = 9.8$ . Note that the value of  $\text{pK}_A$  matches that found from both redox titrations (3–5) and electron transfer (7).

The interaction energy of  $Q_A^-$  with a nearby proton depends on  $k_{\text{obs}}^{H^+}$  and  $k_{\text{obs}}^0$  (see Eq. 1). The energy splitting  $\delta G^0$  is given by (Eq. 1 with  $T = 21^\circ\text{C}$ ):

$$\delta G^0 = k_b T \ln \frac{k_{\text{obs}}^0}{k_{\text{obs}}^{H^+}} = 22 \text{ meV}. \quad (5)$$

The magnitude of  $\delta G^0$  is approximately two orders of magnitude smaller than the optical transition energy of

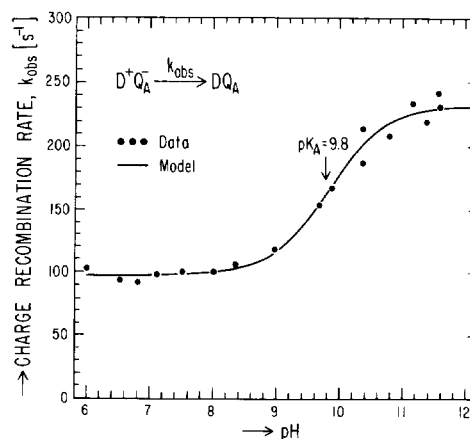


FIGURE 3 The pH dependence of the charge recombination rate  $k_{\text{obs}}$ . The solid line (Model) was calculated using Eqs. 2 and 3 with  $\text{pK}_A = 9.8$ ,  $k_{\text{obs}}^{H^+} = 97 \text{ s}^{-1}$  and  $k_{\text{obs}}^0 = 230 \text{ s}^{-1}$ . Conditions as in Fig. 2, except for varying buffers and pH.

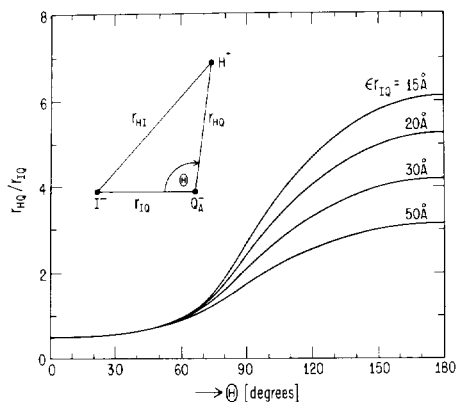


FIGURE 4 The normalized proton to  $Q_A^-$  distance,  $(r_{HQ}/r_{IQ})$ , as a function of the angle,  $\theta$ , between the  $I^- - Q_A^-$  and the  $H^+ - Q_A^-$  axes, for different values of  $\epsilon r_{IQ}$ . The curves were calculated using Eqs. 5 and 6. Inset shows the geometrical arrangement of the reactants.

$Q_A^-$  (9–11). This may explain why the optical spectrum of  $Q_A^-$  is essentially unaffected by the binding of a proton.<sup>2</sup>

From the value of  $\delta G^0$ , one can make a rough estimate of the location of the proton binding site relative to  $Q_A^-$ . We assume that the interaction of the proton with both  $Q_A^-$  and  $I^-$  is electrostatic in origin, as has been done to explain the effects of point charges in other photosynthetic systems (26–29). The change in energy  $\delta G^0$  is caused by the difference in the distance of the proton to  $Q_A^-$  (i.e.,  $r_{HQ}$ ) compared with the distance to  $I^-$  (i.e.,  $r_{HI}$ ). From Coulomb's law we obtain<sup>3</sup>

$$\delta G^0 = \frac{-e^2}{\epsilon r_{HI}} - \frac{-e^2}{\epsilon r_{HQ}}$$

$$= \frac{-e^2}{\epsilon r_{IQ}} \left[ \frac{1}{[1 + (r_{HQ}/r_{IQ})^2 - 2(r_{HQ}/r_{IQ}) \cos \theta]^{1/2}} - \frac{1}{(r_{HQ}/r_{IQ})} \right], \quad (6)$$

where  $e$  is the electronic charge,  $\epsilon$  is the effective dielectric constant,  $r_{IQ}$  is the distance between the  $I^-$  and  $Q_A^-$  and  $\theta$  is the angle between the  $I^- - Q_A^-$  and  $H^+ - Q_A^-$  axes (see insert in Fig. 4). The distance  $r_{IQ}$  has been estimated from spectroscopic measurements to be 8–12 Å (30,31). The value of  $\epsilon$  is difficult to estimate when one is dealing with distances on the scale of atomic dimensions<sup>4</sup> (see references 32–35). Consequently, we have not assumed a specific value of  $\epsilon$ , but have calculated distances in terms of  $\epsilon r_{IQ}$ . By equating Eqs. 5 and 6, permissible combinations of  $(r_{HQ}/r_{IQ})$  and  $\theta$  were determined for the range of  $15 \text{ \AA} \leq \epsilon r_{IQ} \leq 50 \text{ \AA}$  (Fig. 4). The minimum distance of the proton to  $Q_A^-$

<sup>2</sup>Although  $\delta G^0$  represents the change in energy difference between two states (Fig. 1), the shifts in the individual energies are expected to be also small with respect to the optical transition energy.

<sup>3</sup>We assume that the relative distances between the reactants  $D^+$ ,  $I^-$ ,  $Q_A^-$  do not change upon protonation.

<sup>4</sup>It will be of interest to obtain the value of  $\epsilon$  when the distances will have been determined from crystallographic and/or electron nuclear double resonance measurements.

occurs when the proton is situated about halfway between  $I^-$  and  $O_A^-$ , i.e.,  $\theta = 0$ . At this location,  $r_{HQ} \approx 0.5 r_{IQ} \approx 5 \text{ \AA}$ ; the ratio  $r_{HQ}/r_{IQ}$  is essentially independent of  $\epsilon r_{IQ}$ .

In conclusion, we have shown that a proton is associated with  $Q_A^-$  at a distance  $\geq 5 \text{ \AA}$ . This relatively large distance may explain why the protonation has not been observed in either the optical or EPR spectrum of  $(Q_A^-)H^+$ . The pK that we determined (9.8) is in agreement with the values found from redox (3–5) and electron transfer (7) studies.

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