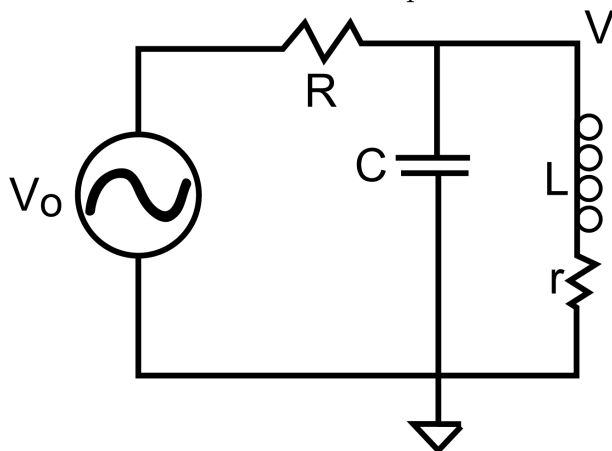


1 RLC resonator with a lossy inductor

The circuit in laboratory exercise 3.1 uses an inductor that has a modest resistance. Let's see how this effects the amplitude of the signal output.



The equivalent impedance of the $(L+r)||C$ circuit is

$$\begin{aligned} Z_{eq} &= \frac{\frac{r+i\omega L}{i\omega C}}{r+i\omega L + \frac{1}{i\omega C}} \\ &= \frac{r+i\omega L}{i+i\omega C(r+i\omega L)} \end{aligned} \quad (1.1)$$

Then

$$\begin{aligned} \frac{V(\omega)}{V_o(\omega)} &= \frac{\frac{r+i\omega L}{i+i\omega C(r+i\omega L)}}{\frac{r+i\omega L}{i+i\omega C(r+i\omega L)} + R} \\ &= \frac{r+i\omega L}{r+i\omega L + R(1+i\omega C(r+i\omega L))} \\ &= \frac{r+i\omega L}{r+R - R\omega^2 LC + i\omega L + i\omega r RC}. \end{aligned} \quad (1.2)$$

We let

$$\omega_0 = \frac{1}{\sqrt{LC}}, \quad (1.3)$$

$$\rho = r/R, \quad (1.4)$$

where the limit of a lossless inductor then corresponds to $\rho \rightarrow 0$, and

$$\tau = L/R. \quad (1.5)$$

In terms of these normalized units,

$$\frac{V(\omega)}{V_o(\omega)} = \frac{\rho + i\omega\tau}{1 + \rho - \left(\frac{\omega}{\omega_0}\right)^2 + i\omega\tau \left(1 + \frac{\rho}{(\omega_0\tau)^2}\right)}. \quad (1.6)$$

The magnitude of the output is then

$$\begin{aligned} \left| \frac{V(\omega)}{V_o(\omega)} \right| &= \sqrt{\frac{\rho^2 + (\omega\tau)^2}{\left(1 + \rho - \left(\frac{\omega}{\omega_0}\right)^2\right)^2 + (\omega\tau)^2 \left(1 + \frac{\rho}{(\omega_0\tau)^2}\right)^2}} \\ &= \sqrt{\frac{\left(\frac{\omega}{\omega_0}\right)^2 + \left(\frac{\rho}{\omega_0\tau}\right)^2}{\left(\frac{1 - \left(\frac{\omega}{\omega_0}\right)^2 + \rho}{\omega\tau}\right)^2 + \left(\frac{\omega}{\omega_0}\right)^2 \left(1 + \frac{\rho}{(\omega_0\tau)^2}\right)^2}} \end{aligned} \quad (1.7)$$

and this has the expected limit as $\rho \rightarrow 0$, i.e.,

$$\left| \frac{V(\omega)}{V_o(\omega)} \right| \xrightarrow{\rho \rightarrow 0} \frac{\frac{\omega}{\omega_0}}{\sqrt{\left(\frac{1 - \left(\frac{\omega}{\omega_0}\right)^2}{\omega\tau}\right)^2 + \left(\frac{\omega}{\omega_0}\right)^2}} \quad (1.8)$$

We return to the general expression of the magnitude of the output and note that for $\omega = \omega_0$, the magnitude of the output is

$$\left| \frac{V(\omega)}{V_o(\omega)} \right| \xrightarrow{\omega = \omega_0} \sqrt{\frac{1 + \left(\frac{\rho}{\omega_0\tau}\right)^2}{\left(\frac{\rho}{\omega_0\tau}\right)^2 + \left(1 + \frac{1}{\rho} \left(\frac{\rho}{\omega_0\tau}\right)^2\right)^2}}. \quad (1.9)$$

The dependence on $\frac{\rho}{\omega_0\tau} = r\sqrt{\frac{C}{L}}$ shows that loss in the inductors is minimized at higher frequencies.

For the components in laboratory 3.1, i.e., $L = 10$ mH, $R = 100$ k Ω , $C = 0.01$ μ F and r measured to be 150 Ω for one inductor, we find $\rho = 0.015$, $\omega\tau = 0.01$, and the decrement in amplitude is large, a factor of 16. This occurs even through ρ is rather small.