

## 4 Electrodiffusion

We previously discussed how the motion of freely dissolved ions and macromolecules is governed by diffusion, the random motion of molecules that results from fluctuating forces. This picture largely holds in the presence of weak forces that merely add a slight bias to the motion, such as the case of electric fields across membrane pores. In particular, in the presence of an electric field the motion is limited by the collisions so that the velocity, as opposed to acceleration, is proportional to the force. We have

$$\begin{aligned}\vec{v}_D(\vec{r}, t) &= \mu \vec{E}(\vec{r}, t) \\ &= -\mu \vec{\nabla} V(\vec{r}, t)\end{aligned}\tag{4.1}$$

where  $\vec{v}_D(\vec{r}, t)$  is known as the drift velocity and  $\mu$  is the mobility. The joke is that life is diffusion with a small drift velocity.

How big are these terms? Without proof, we note that  $|\vec{v}_D| \approx 10^1 \frac{cm}{s} = 10^2 \frac{\mu m}{ms}$  in membranes and  $|\vec{v}_D| \approx 10^{-5} \frac{cm}{s} = 10^{-4} \frac{\mu m}{ms}$  in cytoplasm.

We can now calculate the flux due to the electric field as

$$\begin{aligned}\vec{J}_D(\vec{r}, t) &= C(\vec{r}, t) \vec{v}_D(\vec{r}, t) \\ &= \mu C(\vec{r}, t) \vec{E}(\vec{r}, t) \\ &= -\mu C(\vec{r}, t) \vec{\nabla} V(\vec{r}, t)\end{aligned}\tag{4.2}$$

Thus the total flux, thermal as well as force driven, is

$$\vec{J}(\vec{r}, t) = -D \nabla C(\vec{r}, t) - \mu C(\vec{r}, t) \vec{\nabla} V(\vec{r}, t)\tag{4.3}$$

At equilibrium,  $\vec{J}(\vec{r}, t) = 0$  and thus

$$C(\vec{r}) = C(\vec{r}') e^{-\frac{\mu}{D}(V(\vec{r}) - V(\vec{r}'))}\tag{4.4}$$

or

$$\Delta V = V(\vec{r}) - V(\vec{r}') = -\frac{D}{\mu} \ln \left( \frac{C(\vec{r})}{C(\vec{r}')} \right)\tag{4.5}$$

but we previously showed that this equilibrium potential is just given by the Nernst formula, i.e.,

$$\Delta V = -\frac{k_B T}{e} \ln \left( \frac{C(\vec{r})}{C(\vec{r}')} \right)\tag{4.6}$$

Thus

$$\mu = D \frac{e}{k_B T} \quad (4.7)$$

We can estimate the size of  $|\vec{v}_D|$  for membranes and compare it with the RMS thermal velocity,  $|\vec{v}_{th}|$ , since the concept of the drift velocity holds only if  $|\vec{v}_D| \ll |\vec{v}_{th}|$ . For a membrane, the electric field is of order  $|\vec{E}| = \frac{V}{L} \sim \frac{k_B T}{e} \frac{1}{L}$  so that  $|\vec{v}_D| = \mu |\vec{E}| \sim D \frac{e}{k_B T} \frac{k_B T}{e} \frac{1}{L} \approx \frac{D}{L}$ . On the other hand, we discussed earlier that the scale of the thermal velocity is  $|\vec{v}_{th}| \sim \frac{\lambda}{\tau} \approx \frac{D}{\lambda}$ . Thus

$$|\vec{v}_D| \sim \frac{\lambda}{L} |\vec{v}_{th}| \sim 0.1 |\vec{v}_{th}| \quad (4.8)$$

so that the thermal velocity is relatively small for distances larger than the collision length, a rather intuitive result. On the other hand, for large electric fields may anticipate a deviation from a simple linear relationship between  $|\vec{E}|$  and  $|\vec{v}_D|$ , a topic we will return to.

Since we previously argued that the thermal velocity can be estimated from equipartition, i.e.,  $|\vec{v}_{th}| \sim \sqrt{\frac{k_B T}{m}}$ , we have

$$|\vec{v}_D| \sim c \sqrt{\frac{k_B T}{m c^2}} \frac{\lambda}{L} \quad (4.9)$$

from which we can estimate the absolute size of the drift velocity for membranes to be  $|\vec{v}_{th}| \sim 10 \text{ cm/s}$ .

We can now put all of the formalism together to get a modified flux for the diffusion equation, i.e.,

$$\vec{J}(\vec{r}, t) = -D \left( \nabla C(\vec{r}, t) + \frac{e}{k_B T} C(\vec{r}, t) \vec{\nabla} V(\vec{r}, t) \right) \quad (4.10)$$

At this point, let's cool it with the general form for 3-dimensions and focus on the case of current through a pore of cross sectional area  $A$  that spans a membrane of thickness  $L$ . We further assume that the electric field is uniform (probably not true, but it allows us to make some useful progress), so that  $V(x) = V(0) \frac{x}{L}$ , we have an equation for the electrical current,  $I$ ,

$$I = -e J(x) A = e D A \left( \frac{dC(x)}{dx} + \frac{e}{k_B T} C(x) \frac{\Delta V}{L} \right) \quad (4.11)$$

where  $A$  is the area of the membrane. This equation is in the form of  $\frac{dC(x)}{dx} + \text{const} \cdot C(x) = \text{const}'$ , which we can solve directly to obtain

$$I = e \frac{D}{L} \frac{e V}{k_B T} A \frac{C(L) - C(0) e^{-\frac{eV}{k_B T}}}{1 - e^{-\frac{eV}{k_B T}}} \quad (4.12)$$

This is known as the Nernst-Planck relation. The interesting, and essential feature, is that the  $I - V$  curve is nonlinear for voltage changes on the order of  $\frac{k_B T}{e} \approx 25 \text{ mV}$  away from the reversal potential.

## FIGURE - chapt-4-nernst-plank.eps

If we include the possibility of a valence,  $z = \pm 1, \pm 2, \pm 3$ , etc, the general form of the Nernst-Plank relation becomes

$$I(V) = z^2 e \frac{D}{L} \frac{eV}{k_B T} A \frac{C(L) - C(0) e^{-\frac{zeV}{k_B T}}}{1 - e^{-\frac{zeV}{k_B T}}} \quad (4.13)$$

and is particularly strong for divalents, such as  $Ca^{+2}$ .

The essential physics is that it is "easier" for currents to flow from high concentrations to low concentrations, so that the conductance (slope of the  $I - V$ ) is larger when ions move from high to low, rather than from low to high, concentration. This property is known as rectification, and is a "normal" property of any cell membrane. The limiting currents are Ohmic, with

$$I(V) \rightarrow \begin{cases} z^2 e \frac{D}{L} \left( \frac{eV}{k_B T} \right) AC(L) & \text{if } V \gg \frac{k_B T}{e} \\ z^2 e \frac{D}{L} \left( \frac{eV}{k_B T} \right) AC(0) & \text{if } V \ll -\frac{k_B T}{e} \end{cases}$$

These asymptotic relations are often good approximations, and for the better or worse most physiologists assume that the I-V relation is Ohmic. But do this at your peril. Hagiwara made this assumption about the "Inward Rectifier" for  $K^+$  and came to some erroneous conclusions about his data!

### 4.1 The resting potential of a neuron

In steady-state, the total current across the cell membrane is zero. When only a single ionic species can pass across the membrane, i.e., when the current is carried by only a single ionic species, the steady-state potential is also the equilibrium potential and we recover the Nernst potential. In general, however, multiple ions contribute to the current flow across the membrane. We focus on a cell with just 3 of these,  $Na^+$ ,  $K^+$  and  $Cl^-$ . There is a Nernst-Plank current associated with each ion, i.e.,

$$I_{Na^+}(V) = eA \left( \frac{D}{L} \right)_{Na^+} \frac{eV}{k_B T} \frac{[Na^+]_{in} - [Na^+]_{out} e^{-\frac{eV}{k_B T}}}{1 - e^{-\frac{eV}{k_B T}}} \quad (4.14)$$

$$I_{K^+}(V) = eA \left( \frac{D}{L} \right)_{K^+} \frac{eV}{k_B T} \frac{[K^+]_{in} - [K^+]_{out} e^{-\frac{eV}{k_B T}}}{1 - e^{-\frac{eV}{k_B T}}} \quad (4.15)$$

$$I_{Cl^-}(V) = eA \left( \frac{D}{L} \right)_{Cl^-} \frac{eV}{k_B T} \frac{[Cl^-]_{in} - [Cl^-]_{out} e^{\frac{eV}{k_B T}}}{1 - e^{\frac{eV}{k_B T}}} \quad (4.16)$$

In steady state, we must have

$$I_{Na^+}(V) + I_{K^+}(V) + I_{Cl^-}(V) = 0 \quad (4.17)$$

This condition is satisfied for only a single voltage. The algebra is easy to do if you recall that one should solve for  $e^{\frac{eV}{k_B T}}$  rather than directly for  $V$  (it is also easy if one ignores divalents, like  $Ca^{+2}$ !). The steady state potential is given by (Goldman-Hodgkin-Katz equation)

$$V_{SS} = \frac{k_B T}{e} \ln \frac{\left(\frac{D}{L}\right)_{K^+} [K^+]_{out} + \left(\frac{D}{L}\right)_{Na^+} [Na^+]_{out} + \left(\frac{D}{L}\right)_{Cl^-} [Cl^-]_{in}}{\left(\frac{D}{L}\right)_{K^+} [K^+]_{in} + \left(\frac{D}{L}\right)_{Na^+} [Na^+]_{in} + \left(\frac{D}{L}\right)_{Cl^-} [Cl^-]_{out}} \quad (4.18)$$

The fraction  $\frac{D}{L}$  is often called the permeability, denoted  $P$ . Actually, the permeability is defined with the addition of a mystery fudge factor in front! We will see below that the permeability is dependent on concentration such that the channels are significantly saturated at physiological ion concentrations.

As we mentioned on the first day,  $V_{SS}$  typically has a value of about -50 mV. Ion substitution experiments that maintain a constant total ion concentration have confirmed this relation.

**FIGURE - chapt-4-goldman.doc**

## 4.2 Relations among "Diffusive" processes

Since in the linear limit (i.e.,  $\frac{zeV}{k_B T} \rightarrow \pm\infty$ ) the  $I - V$  relation must reduce to Ohm's law, i.e.,  $I = g \frac{\pi a^2}{L} V$ , we have the relation

$$g = \frac{(ze)^2 DC(L)}{k_B T} \quad (4.19)$$

In general, we recall that the diffusion constant ( $D$ ), the mobility ( $\mu$ ), the viscosity ( $\eta$ ), and the electrical conductance ( $g$ ) are all related. We recapitulate this below,

$$D = g \frac{k_B T}{(ze)^2 C_{ions}} = \mu \frac{k_B T}{ze} = \frac{1}{\eta} \frac{k_B T}{6\pi a} \quad (4.20)$$

where the viscosity relation for a particle of radius  $a$  is not really useful for ions per se.

Lastly, as a practical issue, to convert from "physics" units to "chemical" units, one replaces  $k_B$  by  $R = N_A \cdot k_B$  and  $e$  by  $F = N_A \cdot e$  and measures concentration in molar (moles/liter) rather than ions per  $cm^3$ .

## 4.3 Channel Saturation

One expects that the current will not increase without bound as the ionic concentration increases. This may occur as the kinetics of channels limit the current. A simple example of kinetic limitations is to consider a channel with a single ion binding site. The current will follow saturation characteristics for enzyme binding

(Michaelis-Menten kinetics), so that, e.g., the above limiting case for the current ( $V \gg \frac{k_B T}{e}$ ) becomes

$$I(V) \rightarrow z^2 e \frac{D}{L} \left( \frac{eV}{k_B T} \right) \pi a^2 \frac{C(L)}{1 + \frac{C(L)}{K_{eq}}} \quad (4.21)$$

### FIGURE - chapt-4-channel-conductances.eps

Beyond this effect, steric hindrance and electrostatic repulsion will further limit the currents. Hille has a nice discussion of this.

The bottom line is that the saturation effects are substantial near physiological concentrations. Thus the permeability of channels will be implicitly concentration dependent, and formulas like the Goldman-Hodgkin-Katz equation for the steady-state membrane potential must be applied with permeabilities chosen for, or corrected for, the concentrations used in the particular experiment.

One further expects that the current will not increase without bound as the voltage increases. This may occur as a depletion region develops at one end of the channel, and a space charge region develops at the other end. In effect, the voltage causes a concentration-based limitation. An easy way to estimate this effect is to consider the diffusion equation for a point source (pore opening).

At steady state,  $\frac{\partial C(\vec{r}, t)}{\partial t} = 0$  so that

$$-D \nabla^2 C(\vec{r}) = I(\infty) \delta(\vec{r} - \vec{r}') \quad (4.22)$$

where  $I(\infty) = I(V \rightarrow \infty)$ . This is just Poisson's equation for a point charge, so that

$$-D \nabla^2 C(\vec{r}) = I(\infty) \left( -\frac{1}{4\pi} \nabla^2 \left( \frac{1}{|\vec{r} - \vec{r}'|} \right) \right) \quad (4.23)$$

Thus, integrating over a hemisphere (so that the source doubles compared to the full spherical solution) we have

$$C(\vec{r}) = \frac{I(\infty)}{2\pi D} \frac{1}{r} + \text{Constant} \quad (4.24)$$

This modifies the concentrations on opposite sides of the pore to read

$$C(0) = C_o(0) + \frac{I(\infty)}{2\pi D a} \quad (4.25)$$

$$C(L) = C_o(L) - \frac{I(\infty)}{2\pi D a} \quad (4.26)$$

The current will saturate at a maximal value when the concentrations on opposite sides of the pore are equal, i.e.,  $C(0) = C(L)$ . Then the electrical current is

$$eI(\infty) = e\pi D a [C_o(L) - C_o(0)] \quad (4.27)$$

An order of magnitude estimate of this space charge effect for a channel with a  $10\text{\AA}$  pore diameter is  $30 \times 10^{-12}$  Amperes. This is close to the measured saturation levels for  $K^+$  flow through the  $Ca^{2+}$ -activated  $K^+$  channel.

**FIGURE - chapt-4-space-charge-?.eps**

Lastly, we can substitute the corrected form for the concentration back into the Nernst-Planck relation. Thus, for example, the above limiting case for the current through the pore becomes ( $V \gg \frac{k_B T}{e}$ ) becomes

$$I(V) \rightarrow \frac{z^2 e \frac{D}{L} \left( \frac{eV}{k_B T} \right) \pi a^2 C(L)}{1 + z^2 e \frac{a}{2L} \left( \frac{eV}{k_B T} \right)} \quad (4.28)$$