

Receptive Field

$$z(t) = g \left[I_0 + \int d^2r \int_{-\infty}^t dt' R(\vec{r}, t-t') S(\vec{r}, t') \right]$$

\uparrow Prob of spiking or instantaneous spike rate
 \uparrow Receptive Field
 \uparrow Stimulus

when the stimulus driven part is small compared to the background z_0

$$z(t) = g(I_0) + \frac{dg}{dI} \Big|_{I=I_0} \int d^2r \int_{-\infty}^t dt' R(\vec{r}, t-t') S(\vec{r}, t')$$

\uparrow $\equiv z_0$
 \uparrow $\equiv g'$

Probability of no spikes in the interval $[0, t]$ and one spike in $(t, t+dt)$ is

$$e^{-\int_0^t dt' z(t')} \cdot z(t)$$

for a Poisson process

Modulation in firing rate by convolution of stimulus with receptive field

For $R(\vec{r}, t) \equiv \sum_n \lambda_n \underbrace{F_n(\vec{r}) G_n(t)}_{\text{modes of receptive field}}$

$$z(t) = z_0 + g' \sum_n \lambda_n \int d^2r F(\vec{r}) \int_{-\infty}^t dt' G_n(t-t') S(\vec{r}, t')$$

For $S(\vec{r}, t) = X(\vec{r}) \delta(t)$ (In practice rapidly changing stimulus)

$$z(t) = z_0 + g' \sum_n \lambda_n G_n(t) \int d^2r F(\vec{r}) X(\vec{r})$$

Rate depends on spatial pattern

Let us reverse process and ask if we can reconstruct a stimulus from the spike train

For simplicity, let us ignore space. The previous description of receptive field gives

$$\rightarrow Z(t) = z_0 + g' \int_{-\infty}^t dt' R(t-t') S(t')$$

Averaged and smoothed
spike train

↓ Fourier Transform

$$\tilde{Z}(\omega) = z_0 \delta(\omega) + g' \tilde{R}(\omega) \tilde{S}(\omega)$$

$$\tilde{S}(\omega) = \frac{\tilde{Z}(\omega) - z_0 \delta(\omega)}{g' \tilde{R}(\omega)} \quad \text{not very informative}$$

Optimal Reconstruction

How well can we reconstruct stimulus from a given spike train

$$\text{Let } \Lambda_x(t) \equiv \sum_{\text{spikes}} \delta(t-t_s) \quad \begin{array}{l} \text{spike train} \\ \text{for } x\text{-trial} \end{array}$$

$$\rightarrow \text{predicted } S_x(t) = \int_{-\infty}^t dt' T(t-t') \Lambda_x(t')$$

↑
Transfer Function ($\tilde{T}(\omega)$ & $\tilde{R}(\omega)$?)

$$\text{Compare } S_x^{\text{actual}}(t) \text{ vs. } S_x^{\text{predict}}(t)$$

$$\text{Integrated Error} \equiv \sum_x \int dt [S_x^{\text{actual}}(t) - S_x^{\text{predict}}(t)]^2$$

↓ Fourier Transform
(different ω 's are uncorrelated)

$$\text{Error} = \sum_{\omega} \left[\tilde{S}_{\alpha}^{\text{pred}}(\omega) - \tilde{S}_{\alpha}^{\text{actual}}(\omega) \right]^2$$

$$= \sum_{\omega} \left| \tilde{T}(\omega) \tilde{\Lambda}_{\alpha}(\omega) - \tilde{S}_{\alpha}^{\text{actual}}(\omega) \right|^2$$

$$= \sum_{\omega} \left[\tilde{T}(\omega) \tilde{T}^*(\omega) \tilde{\Lambda}_{\alpha}(\omega) \tilde{\Lambda}_{\alpha}^*(\omega) + \tilde{S}_{\alpha}^{\text{act}}(\omega) \tilde{S}_{\alpha}^{\text{act}*}(\omega) - \tilde{T}(\omega) \tilde{\Lambda}_{\alpha}(\omega) \tilde{S}_{\alpha}^{\text{act}*}(\omega) - \tilde{T}^*(\omega) \tilde{\Lambda}_{\alpha}^*(\omega) \tilde{S}_{\alpha}^{\text{act}}(\omega) \right]$$

$$\frac{\partial \mathcal{E}}{\partial \tilde{T}^*(\omega)} = \sum_{\omega} \left[\tilde{T}(\omega) \tilde{\Lambda}_{\alpha}(\omega) \tilde{\Lambda}_{\alpha}^*(\omega) - \tilde{\Lambda}_{\alpha}^*(\omega) \tilde{S}_{\alpha}^{\text{act}}(\omega) \right]$$

holds for each value of $\tilde{T}^*(\omega)$, so really differentiation w.r.t. a scalar.

$$\text{Let } \frac{\partial \mathcal{E}}{\partial \tilde{T}^*(\omega)} = 0 \quad \therefore \tilde{T}(\omega) |\tilde{\Lambda}_{\alpha}(\omega)|^2 = \tilde{\Lambda}_{\alpha}^*(\omega) \tilde{S}_{\alpha}^{\text{act}}(\omega)$$

$$\therefore \tilde{T}(\omega) = \frac{\sum_{\omega} \tilde{\Lambda}_{\alpha}^*(\omega) \tilde{S}_{\alpha}^{\text{act}}(\omega)}{\sum_{\omega} |\tilde{\Lambda}_{\alpha}(\omega)|^2}$$

Cross Power
Power in Spike Train

$$\text{And } \tilde{S}_{\alpha}^{\text{pred}}(\omega) = \tilde{T}(\omega) \tilde{\Lambda}_{\alpha}(\omega)$$

Best Filter

The transfer function is just the cross-correlation, divided by the autocorrelation to correct for trends of correlation among the neuronal response $\Lambda(\omega)$.

We shall see next how the filter has a simple form - and reduces to the expression for the reverse correlation or spike-triggered average, when the correlation between spikes is zero. This holds for low spike rates; at high rates the refractory period leads to correlations in the rate.

How do we interpret the result for the optimal filter? Look at the case of $\Lambda(\omega)$ for a spike train

First - the numerator - is just the cross power, or the Fourier Transform of the cross-correlation, i.e.

$$\sum_{\sigma} \tilde{\Lambda}_{\sigma}(\omega) \tilde{S}_{\sigma}(\omega) = \sum_{\sigma} \int_{-\infty}^{\infty} d\tau e^{-i\omega\tau} \int dt \Lambda_{\sigma}(t) S_{\sigma}(t-\tau)$$

$$\text{But } \sum_{\sigma} \Lambda(t) \rightarrow \sum_{\text{spike times}}^N \delta(t-t_s)$$

$$\therefore \sum_{\sigma} \tilde{\Lambda}_{\sigma}(\omega) \tilde{S}_{\sigma}(\omega) = \int_{-\infty}^{\infty} d\tau e^{-i\omega\tau} \sum_{\substack{\text{spike} \\ \text{times}}}^N \int dt \delta(t-t_s) S(t-\tau)$$

$$= \int_{-\infty}^{\infty} d\tau e^{i\omega\tau} \sum_{\substack{\text{spike} \\ \text{times}}}^N S(t_s - \tau)$$

Spike Trigger
Stimulus Avg.

Second - the denominator is just the power, or the Fourier transform of the auto-correlation

$$\sum_{\sigma} |\tilde{\Lambda}_{\sigma}(\omega)|^2 = \int_{-\infty}^{\infty} d\tau e^{i\omega\tau} \sum_{\substack{\text{spike} \\ \text{times}}}^N \int dt \delta(t-t_s) \delta(t-t_s+\tau)$$

$\delta(t_s - t_s + \tau) = \delta(\tau)$

$N \delta(\tau)$

$$= N$$

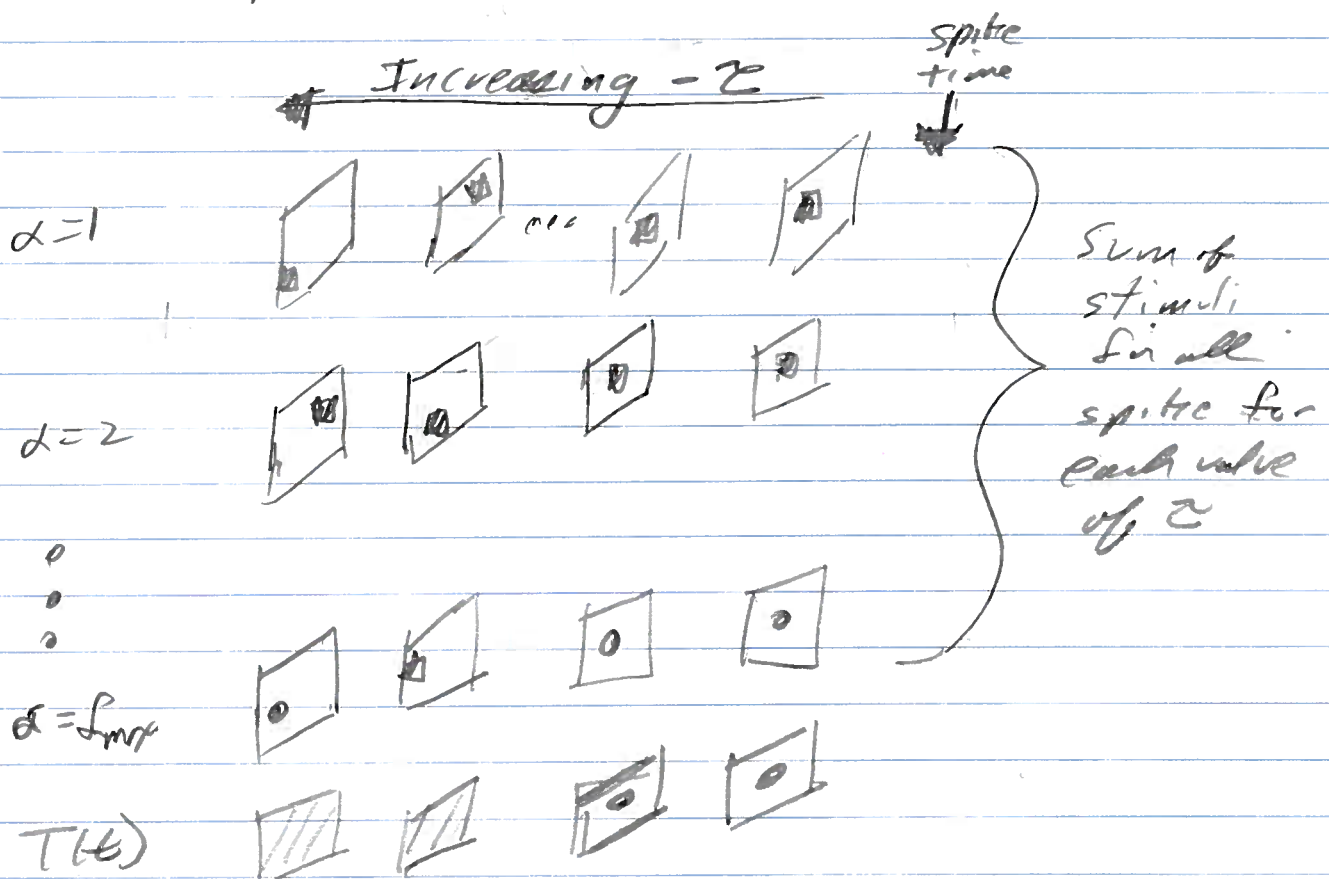
∴ We can write that for the case of uncorrelated spikes

$$T(\tau) = \int_{-\infty}^{\infty} d\alpha e^{i\omega\tau} \tilde{T}(\alpha) = \frac{1}{N} \int_{-\infty}^{\infty} d\alpha e^{i\omega\tau} \sum_{\alpha} \tilde{R}_{\alpha}(\alpha) \tilde{S}_{\alpha}(\alpha)$$

$$= \frac{1}{N} \sum_{\text{Spike time}} S(t_s - \tau)$$

Reverse Correlation Formula.

This formula is very easy to apply. Wait for a spike, sum up stimulus at time τ is the past, for all



Notes on Singular Value Decomposition

$$R(r, t) \equiv \sum_n \lambda_n F_n(r) G_n(t)$$

$$\left. \begin{array}{l} \text{with } \int d^2r F_n(r) F_m(r) = \delta_{nm} \\ \text{and } \int dt G_n(t) G_m(t) = \delta_{nm} \end{array} \right\} \text{Orthogonal} \\ \text{Functions}$$

Consider contraction to symmetric correlation matrix

$$C(t, t') \equiv \int d^2r R(r, t) R(r, t')$$

$$= \sum_n \sum_m \lambda_n \lambda_m \underbrace{\int d^2r F_n(r) F_m(r)}_{\delta_{nm}} G_n(t) G_m(t')$$

$$= \sum_n \lambda_n^2 G_n(t) G_n(t')$$

$$\text{Then } \int dt' C(t, t') G_m(t') = \sum_n \lambda_n^2 G_n(t) \underbrace{\int dt' G_n(t') G_m(t')}_{\delta_{nm}}$$

$$\int dt' C(t, t') G_m(t') = \lambda_n^2 G_m(t) \quad \leftarrow \text{Eigenvalue Problem}$$

and the $F_n(t)$ are found by

$$\int dt R(r, t) G_m(t) = \sum_n F_n(r) \underbrace{\int dt G_n(t) G_m(t)}_{\delta_{nm}}$$

$$= F_m(r)$$

Bottom line: Space-time modes from measured RF